

# Supporting information for Design of a fiber-cavity ion-trap for a high efficiency and high rate quantum network node

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## Graphical abstract

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## Coupling between ion and fiber cavity

A common ion-cavity system can be described by the parameters  $(g, \kappa, \gamma)$ , where  $g$  is the coherent interaction rate between the particle and cavity,  $\kappa$  is the HWHM (half-width half-maximum) of the cavity transmission line (intracavity field decay at a rate  $2\kappa$ ), and  $\gamma$  is the FWHM (full-width half-maximum) of the spontaneous line (spontaneous decay at a rate  $\gamma$ ).

Note that

$$\kappa = \frac{2\pi\delta\nu}{2} \approx \frac{c[2 - (1 - \mathcal{L}_c)(\mathcal{R}_l + \mathcal{R}_r)]}{4n_c L}, \quad (1)$$

where  $\delta\nu$  is the FWHM frequency of the cavity transmission line,  $\mathcal{L}_c$  is the single-pass effective cavity loss,  $L$  is the cavity length,  $n_c$  is the refractive index of the cavity, and  $\mathcal{R}_l$  and  $\mathcal{R}_r$  are the reflectivity of the left and right cavity mirrors, respectively. In this equation,  $[2 - (1 - \mathcal{L}_c)(\mathcal{R}_l + \mathcal{R}_r)]$  represents an approximate round-trip loss, and this approximation only holds when the cavity loss  $\mathcal{L}_c$  is very small and the reflectivity of the mirror  $\mathcal{R}_l$  or  $\mathcal{R}_r$  is very large.

It is worth noting that the intensity of a single reflection is affected by the clipping loss because the size of the cavity surface of the fiber cavity is finite. The clipping loss is defined as

$$\mathcal{L}_{cl} = \exp\left(-\frac{D^2}{2w_m^2}\right), \quad (2)$$

where  $w_m$  is the radius of the cavity mode at the position of the fiber surface, and  $D$  is the diameter of the effective reflective surface.

The coupling coefficient  $g$  is evaluated by

$$g = \frac{\mu_d E}{\hbar}, \quad (3)$$

where  $\mu_d$  is the electric dipole moment of the particle and  $E$  is the electric field strength of a single photon at the position of the particle. In general, the single-photon electric field intensity is

$$E = \sqrt{\frac{\hbar\omega}{2\epsilon_0 V}} \psi(\vec{r}_a), \quad (4)$$

where  $V = \pi w_0^2 L/4$  is the modal volume,  $w_0$  is the waist radius,  $\vec{r}_a$  is the position of the ion and  $\psi(\vec{r})$  is the normalized cavity mode distribution where  $\max_{\vec{r}} |\psi(\vec{r})| = 1$ . Using the Gaussian mode to approximate the cavity mode, the size of the beam waist satisfies  $w_0^2 = z_0 \lambda / \pi$ , where  $z_0$  is the Rayleigh length and  $\lambda$  is the wavelength.

Define the single-atom cooperativity parameter  $C_1 = g^2 / \kappa \gamma$ . Then, the quantum efficiency of single-photon sources  $\eta_q$  in the bad cavity regime ( $\kappa \gg g^2 / \kappa \gg \gamma$ ) is given by

$$\eta_q = \frac{2C_1}{2C_1 + 1}. \quad (5)$$

In the strong regime ( $g \gg \kappa, \gamma$ ), the quantum efficiency of single-photon sources  $\eta_q$  is given by

$$\eta_q = \frac{2C_1}{2C_1 + 1} \frac{2\kappa}{2\kappa + \gamma}. \quad (6)$$

By changing the structure of the cavity, we can change the values of  $g$  and  $\kappa$  and thus achieve Purcell enhancement.

$\kappa$  in the above equation contains many additional transmission channels and loss channels, while the available single photons are only transmitted in a specific transmission channel. Therefore, it is necessary to consider the percentage of transmission in the total channels, that is,

$$\eta_t = \frac{\mathcal{T}_o}{2 - (1 - \mathcal{L}_c)(\mathcal{R}_l + \mathcal{R}_r)}, \quad (7)$$

where  $\mathcal{T}_o$  is the transmissivity of the output channel.

In addition to  $\eta_q$ , the matching between cavity modes and fiber modes also limits the overall coupling efficiency. This matching efficiency  $\eta_{fc}$  is given by

$$\eta_{fc} = \frac{4}{\left(\frac{w_{f,o}}{w_{m,o}} + \frac{w_{m,o}}{w_{f,o}}\right)^2 + \left(\frac{\pi n_{f,o} w_{f,o} w_{m,o}}{\lambda R_o}\right)^2}, \quad (8)$$

where  $n_{f,o}$  is the refractive index of the output fiber,  $w_{f,o}$  is the waist radius of the corresponding fiber mode,  $w_{m,o}$  is the radius of the cavity mode at the position of the mirror of the output fiber and  $R_o$  is the radius of curvature of this mirror.

In this system, the tunable parameters are the cavity length  $L$ , coating parameters  $(\mathcal{R}, \mathcal{T}, \mathcal{L})$ , structure of the cavity, and type of fiber.